

Laminar-Turbulent Transition on a Cambered NACA 16-009 Airfoil at Low Speed

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Abstract – Infrared thermography, force measurements, and oil flow visualizations are used to investigate the flow patterns around a cambered NACA16-409 airfoil at low Mach number. This cambered profile is widely used for propellers despite the lack of knowledge concerning its flow characteristics. The post-processing of thermograms relies on the analysis of the surface temperature gradient and identification of inflexion points in the temperature distribution. The observations made on the thermograms, based on the distribution of the temperature and Stanton number, are substantiated by the oil flow visualizations. RANS simulations with a transitional SST $k - \omega \& \gamma$ turbulence model corroborate the analysis and deliver detailed insight in the flow around the airfoil. Depending on the angle of attack, three distinct flow patterns have been identified: laminar flow with early separation, laminar separation bubble with trailing edge separation, and turbulent flow with trailing edge separation. The shift between the last two regimes occurs sharply. The prediction capability of the transitional RANS simulations and XFOIL in terms of separation as well as reattachment location are compared with the experimental results. The force-coefficients dependency on the angle-of-attack, obtained by experiments, XFOIL, and RANS simulations, bear the traces from these flow patterns. Copyright © 2020 The Authors. Published by Praise Worthy Prize S.r.l. This article is open access published under the CC BY-NC-ND license (http://creativecommons.org/licenses/by-nc-nd/3.0/).

Keywords: Boundary Layer,	Laminar-Turbulent	Transition,	Separation,	Laminar	Separation
Bubble, Thermogra	aphy, NACA16-Serie	S			

	Nomenclature	V
C	Heat conceity [1/kg/K]	x
C	Durance (finite to the test of tes	α
c_D	Drag coefficient	ε
\mathcal{L}_{f}	Local skin friction coefficient	ρ
C_L	Lift coefficient	∞
C_{ld}	Design lift coefficient	r
C_p	Pressure coefficient	W
С	Chord [m]	~
c_p	Isobar heat capacity [J/kg/K]	CFD
e_a	Approximate relative error	GCI
e_{ext}	Extrapolated relative error	IR
F_i	Force component [N]	NACA
h	Convective heat transfer coefficient	
	$[W/m^2/K]$	LSB
k	Heat conductivity [W/m/K]	RANS
k_{ii}	Interpolation coefficient	TE
l_{sep}	Separation length [m]	
Ma	Mach number	
Pr	Prandtl number	
<i>q</i> _{conv}	Convective heat flux [W/m ²]	The
Re _c	Reynolds number based on chord	for us
$\operatorname{Re}_{\theta,t}$	Reynolds number based on transitional	found
- ,-	momentum thickness	series
St	Stanton number	delav
Т	Temperature [K]	and a
U_i	Channel Voltage [V]	high l

V	Velocity [m/s]
x	Coordinate
α	Angle of attack [°]
ε	Emissivity
ρ	Mass density [kg/m ³]
∞	Freestream value
r	Recovery value
W	Wall value
~	Non-dimensionalized value
CFD	Computational Fluid Dynamics
GCI	Grid Convergence Index
IR	InfraRed
NACA	National Advisory Committee for
	Aeronautics
LSB	Laminar Separation Bubble
RANS	Reynolds-Averaged Navier-Stokes
TE	Trailing Edge

I. Introduction

The NACA 16-series has been specifically designed for use in aircraft propeller blades. Hence, it is widely found in propeller designs as in [1]-[9]. The NACA 16series features flat pressure distributions intended to delay the onset of drag rise due to compressibility issues, and a small leading-edge radius. Both features result in high lift to drag ratios at low angle of attack in cruise

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[10] but also suffer from a reduction in efficiency at high angle of attack due to the separation triggered by the sharp leading edge [1]. Therefore, the investigation of the low-Mach operating conditions, typical of takeoff/landing and the initial climb segment, is important.

Yet, to the authors' knowledge, studies on this particular profile are limited to the aerodynamic performance [11], [12], the data collected in [1], the validation results in [2], which are more than three decades old, and the roughness-induced transition study in the transonic regime for a symmetric NACA 16-series airfoil in [24]. In the study of the flow around immersed bodies such as airfoils, the knowledge of the actual laminar/turbulent/separated state of the boundary layer is of high importance given its implications on the performance and behavior of the body under scrutiny [13]-[15].

The change in skin friction associated with transition or separation results in a change in surface heat transfer [13]. Therefore, infrared thermography as a means of characterizing the status of the boundary layer, has become a popular technique because of its ease of use and its non-intrusiveness [15]-[36]. However, this technique alone is not sufficient to discriminate between some states. A detailed discussion of the basic principles and measurement errors can be found in [34], [35].

Recent developments reported on: (1) the importance of the heating technique, especially when external heating is applied [17], [23], [25], [27], [28]; (2) on the issues selecting the appropriate material for the body [23], [28]-[31], either for the bulk body [28, [29], [31], or by applying a paint or a coating to circumvent the properties of the core at the surface where heat exchanges occur [23], [28]-[30]; (3) and on the choice of a post-processing technique to infer quantitative results.

These techniques range from simple contrast-based algorithms to differential analysis of a sequence of images in space [16], [18]-[20], [26], [33] or time [25].

The present work intends to combine infrared thermography, force measurements, and oil flow visualizations to investigate the flow patterns around a cambered NACA16-409 airfoil ($C_{ld} = 0.3675$) at low Mach number (Ma=0.1, Re_c = 330,000). First, the experimental set-up for the investigation by infrared thermography is described before explaining the consolidating experiments (forces measurements and oil flow visualization). The RANS model is then depicted.

All measurements are then consolidated in the next section before drawing conclusions.

II. Experimental Setup and Theory

II.1. Measurement Principle

The convective heat transfer rate \dot{q}_{conv} is given by Newton's cooling law:

$$\dot{q}_{conv} = h(T_w - T_r) \tag{1}$$

where h is the convective heat transfer coefficient, T_w is

the wall temperature and T_r is the recovery temperature of the flow ($T_r = T_{\infty}$ for low subsonic flows). With ρ_{∞} the mass density, V the velocity, and c_p the isobar heat capacity, the Stanton number relates the heat transferred into a fluid to its thermal capacity:

$$St = \frac{h}{\rho_{\infty} V c_p} \tag{2}$$

For subsonic air flow (for which the Prandtl number is $Pr\approx 1$), the heat and momentum transfer in the boundary layer are related through Reynold's analogy:

$$St \approx \frac{C_f}{2}$$
 (3)

Consequently, the skin friction coefficient and heat transfer rate are directly related through the convective heat transfer coefficient so that an increase in the turbulent behavior of the boundary layer results in an increased heat transfer, thence a lower wall temperature.

II.2. Measurement Setup

The IR radiation emitted from the wing profile is measured with a FLUKE Ti50 camera equipped with a vanadium oxide uncooled sensor. The camera operates in the 8 µm to 14 µm spectral band with a Noise Equivalent Temperature Difference (NETD) of \leq 70 mK at 30 °C and has a resolution of 320×240. Since knowledge of the absolute temperature *T* is not required to gather the information on the laminar/turbulent/separated behavior of the boundary layer, temperature profiles are nondimensionalized with respect to the minimum and maximum temperature of the thermogram at hand:

$$\tilde{T} = \frac{T - T_{\min}}{T_{\max} - T_{\min}}$$
(4)

In several recent contributions, the external heating is provided by a single source [17], [25], [27]-[28] with difficulties resulting from the non-uniform heating [27], [29]. In the present setup, an array of spots is used for external heating of the wing profile prior to the measurements (Fig. 1).

The array consists of HOENLE Superspot 575 solar simulation spots that are arranged to yield a uniformly distributed 800 W/m² irradiance (according to [37]) over a given surface [38].

II.3. Wing Profile

A cambered version of the NACA 16-009 airfoil with $C_{ld} = 0.3675$ is used since it is representative for the blade section at 75%-radius of several propeller families.

The profile is milled from OBOMODULAN boards made of high-density polyurethane foam (PUR). For the present tests, boards of 300 kg/m^3 are used while densities of up to 1600 kg/m^3 can be reached if required.

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Fig. 1 Experimental set-up showing the wing profile with turbulence generator, the IR-camera and its Field Of View (FOV), and the array of spots used in the heating phase

Such high-density boards can be machined to parts with a homogeneous and smooth surface while having sufficient mechanical properties for wind-tunnel testing.

The thermal properties of high-density PUR are summarized in Table I. Although not reported in [29], [31], high-density polyurethane foam is a suitable candidate for infrared thermography, especially in the 300 kg/m³ range, because of its low specific heat capacity *C* and heat conductivity *k*, while having a homogeneous structure. Its high emissivity ε alleviates the issue of infrared reflections [23], [30]. In accordance with established practice [27], [28], [39], the wing profile is equipped with a turbulence generator at the leading edge (Fig. 1) in order to provide the disturbed (turbulent) flow as a witness.

II.4. Process

The experiment starts with the heating phase in which external heating is provided by the array of spots (Fig. 1).

Fig. 2 shows the temperature distribution over the profile's surface obtained with the present arrangement.

The temperature variation over the main surface of interest after the heating phase is of the order of 0.5 K.

The high emissivity of polyurethane foam and the granularity of the surface in terms of IR or visible wavelengths, yield a non-reflecting surface in the infrared range thereby circumventing known issues with other materials [25], [30]. Then, the tunnel wall, which features an access window for the IR-camera, is put in place before measurements are taken once the desired operating Mach number is reached. The measurements can typically be sustained for 10 minutes. The temperature difference between the undisturbed flow of interest and the disturbed flow during the measurement phase is higher than 0.8 K.

All infrared measurements are performed in the closed section of the low-turbulence (Tu $\leq 0.07\%$) wind-tunnel for a Mach number (Ma) of 0.1 and a Reynolds number based on the chord (Re_c) of 330,000.

TABLE I TYPICAL MATERIAL PROPERTIES OF HIGH-DENSITY POLYLIFETHANE FOAMS

	POLYURETHANE FOAMS			
ρ (kg/m ³)	C (J/kg/K)	k (W/m/K)	Е	
300	1120	0.05	0.9	
500	1400-1500	0.05-0.075		
1600	2000-2100	0.12-0.17		

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Fig. 2. Average non-dimensionalized temperature profile (\tilde{T}) on the wing profile and corresponding thermogram measured at the end of the heating phase

III. Consolidating Experiments

III.1. Force Measurements

Force measurements were made in identical flow conditions with an in-house designed and built threecomponent balance shown in Fig. 3. Independent measurement of the lift, drag, and pitching moment with negligible interactions is obtained thanks to pivot-rods (shown in purple and yellow in Fig. 3) that free one single degree of freedom at a time. It uses three Futek LSM400 load cells and IAA300 amplifiers. The Box-Behnken calibration scheme has 18 points and yields the coefficients of the polynomial with interaction terms:

$$F_{i} = k_{i0} + k_{i1}U_{1} + k_{i2}U_{2} + k_{i3}U_{3} + k_{i12}U_{1}U_{2} + k_{i13}U_{1}U_{3} + k_{i23}U_{2}U_{3} (i = 1,2,3)$$
(5)

that define the response surface between the known forces F_i and the measured load cell output voltages U_i .

The measurement accuracy is of 1.5% for lift and 4.6% for drag.



Fig. 3. In-house designed and built 3-components balance

III.2. Oil Flow Visualizations

The oil flow visualizations are performed under flow conditions identical to the infrared ones but in an open test-section wind-tunnel. Optical pictures have been taken after sufficient running time for the proprietary oils and pigments mixture to dry.

IV. Numerical Simulations

IV.1. Spatial Discretization

The reference C-grid consists of 399×336 points respectively on the airfoil surface and in the wall-normal direction. It has an average y + -value of 0.21 on the airfoil surface. Angle-of-attack corrections according to [40] are applied for matching operating conditions but never exceed 0.3°.

A grid convergence study has been completed according to Celik et al. [41] with a 285×240 coarse grid and a 559×470 fine grid. The drag coefficient C_D and non-dimensionalized point of separation location l_{sep}/c at 3° angle of attack are considered as indicators. The indicators from Table II as well as the errors and the Grid Convergence Index (GCI) from Table III suggest that satisfactory grid independence is achieved.

IV.2. Numerical Model

Numerical simulations and oil flow experiments consolidate the observations made on the basis of the thermal images. The 2D steady- RANS simulations of the undisturbed flow around the airfoil in free-field are performed with the transitional SST k- ω and γ turbulence model of ANSYS Fluent which couples the shear-stress transport k-w model developed by Menter [43] with a single transport equation for the intermittency γ and transition momentum thickness Reynolds number [44], [45]. The model has provisions for separation-induced transition ensuring rapid transition once laminar separation occurs. Constant-temperature no-slip wall conditions are applied for the airfoil as a model for the temperature profile measured at the start of the experiment whereas far-field conditions apply on the exterior boundaries.

TABLE II GRID INDEPENDENCE GENERAL RESULTS (J=1.61)

	Experiments	Fine	Regular	Coarse
C_D	0.0089	0.0086	0.0086	0.0088
l_{sep}/c	0.74 0.747 [25]	0.7265	0.7261	0.7273

TABLE III GRID INDEPENDENCE INDICATORS ACCORDING TO CELIK ET AL. [41]

		Fine-Regular	Regular-Coarse
	e_{a} (%)	0.22	2.47
C_D	e _{ext} (%)	0.02	2.40
	GCI (%)	0.02	0.29
	e_{a} (%)	0.06	0.16
l _{sep} / c	<i>e</i> _{ext} (%)	0.03	0.09
	GCI (%)	0.04	0.16

The mass and energy conservation equations and the transport equations are solved using second-order schemes.

V. Measurements and Post-Processing

The combined results from infrared thermography, oil flow visualizations, and CFD are shown in Figs. 4-6 respectively for angles of attack of 0° , 3° , and 9° . Each of these angles illustrates the distinctive flow patterns shown in Fig. 7.

Since the chordwise surface temperature gradient is indicative of the state of the boundary layer [30], [34], we have used the second derivative of the temperature distribution to identify the presence of significant above measurement noise inflexion points which are related to events in the aerodynamic behavior of the boundary layer [20], [23], [24], [35].

For this purpose, a low-pass filter is applied to get rid of spurious inflexion points related to measurement-noise as in [23].

The inflexion points are shown in Figs. 4-6 and denote major events occurring in the boundary layer. The advantage of the inflexion-point method over differential methods is that it does not require a noticeable change in event location between two flow conditions [18], [26].

Indeed, the flow pattern of Figs. 4 for example is maintained between 0° and 3° angle-of-attack. The thermogram at 0° angle-of-attack is shown in Figs. 4 together with the non-dimensionalized temperature profile, the oil-flow visualization, as well as the pressure and skin-friction coefficients, Stanton number profile, and chordwise velocity contours issued from the numerical simulations.

The rows of pixels used in the spanwise averaging process, which constitutes a first low-pass filter [35], and the distinctive flow features are also highlighted. Because of Eq. (3), the decreasing skin friction from the leading edge to 75%-chord goes in pair with a decrease in heat transfer from the hot surface resulting in an increase of the surface temperature and consequently of the intensity in the thermogram.

This is also evidenced by the difference in surface temperature gradient between the undisturbed and the disturbed turbulent flow. At 73%-chord for the experiments (first inflexion point in \tilde{T}) and 76%-chord for the CFD, the boundary layer separates.

The low speed flow along the wall in the separated region heats up leading to high wall-temperatures, thence a small temperature gradient in the wall normal direction that thwarts convective heat transfer with a consequent drop in Stanton number.

Close to 85%-chord, the second inflexion point in \tilde{T} marks where the separated shear flow becomes turbulent.

The location of transition for the CFD results shown in Fig. 7 is educed from the empirical correlation on the transitional momentum thickness Reynolds number $\text{Re}_{\theta,t}$ in [46], which are similar to Abu-Ghannam and Shaw [47].

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Figs. 4. Results at 0° angle of attack (suction side): a) pressure and skin friction coefficients from CFD (*C_p* and *C_f*); b) non-dimensionalized temperature profile *T̃* from IR-thermography with emphasis on the inflection points (∇), and Stanton number distribution from CFD (St);
c) thermogram with indication of the rows used for spanwise averaging (←→) of the undisturbed and disturbed (turbulent) flows; d) oil flow visualization; e) chordwise velocity contours from CFD

Downstream of this point, the Stanton-profile features higher heat transfer to the fluid. The turbulent flow does not reattach as is evidenced from the skin-friction profile, the oil-flow visualization, and velocity profiles.

However, the non-reattachment cannot be deduced from the thermogram which features an additional inflexion point where the temperature profile matches that of the disturbed turbulent witness-flow. This third inflexion point corresponds to increased skin friction in the reversed wall-bounded flow and transfer of additional heat.



Figs. 5. Results at 3° angle of attack (suction side): a) pressure and skin friction coefficients from CFD (*C_p* and *C_f*); b) non-dimensionalized temperature profile *T̃* from IR-thermography with emphasis on the inflection points (∇), and Stanton number distribution from CFD (St);
c) thermogram with indication of the rows used for spanwise averaging (←→) of the undisturbed and disturbed (turbulent) flows; d) oil flow visualization; e) chordwise velocity contours from CFD

Figs. 5 illustrate the flow pattern found when a Laminar Separation Bubble (LSB) occurs, here for 3° angle-of-attack. The laminar boundary layer separates at 74%-chord (first inflexion point in \tilde{T}) with a 1% discrepancy with the CFD. At 85%-chord, the separated shear flow becomes turbulent (second inflexion point in \tilde{T}) within 2%-chord of the CFD prediction. The resulting increased momentum exchange in the direction perpendicular to the wall results in higher skin friction in the separated bubble and the fluid adjacent to the wall is more effectively heated when compared to the Stanton-distribution from Fig. 4(b).

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Figs. 6. Results at 9° angle of attack (suction side): a) pressure and skin friction coefficients from CFD (*C_p* and *C_f*); b) non-dimensionalized temperature profile *T̃* from IR-thermography with emphasis on the inflection points (∇), and Stanton number distribution from CFD (St);
c) thermogram with indication of the rows used for spanwise averaging (←→) of the undisturbed and disturbed (turbulent) flows; d) oil flow visualization; e) chordwise velocity contours from CFD

Eventually, the turbulent separated flow reattaches at 95%-chord (third inflexion point in \tilde{T}) before separation occurs again up-stream of the Trailing Edge (TE). The agreement with the oil-flow visualizations is excellent.

Such agreement was also reported for an SD7037 airfoil at low Reynolds number in [23] and a symmetric NACA 16-series in [24]. However, it should be noted that the possible occurrence of reattachment cannot be deduced from the thermograms or temperature profiles alone as is evidenced from the comparisons of Figs. 4(b) and (c) and 5(b) and (c).



Fig. 7. Separation, reattachment, transition, and critical location (x/c) versus angle of attack α for the suction side of the NACA16009 airfoil (C_{ld} =0.3675)

At 9° angle-of-attack as shown in Figs. 6, the flow over the airfoil is entirely turbulent and differences between the disturbed and undisturbed flow vanish. The comparison of the gradient with Figs. 5 stresses the need for a witness flow since inferring the laminar or turbulent behavior of the flow from the gradient of the undisturbed flow alone only increases the risk for erroneous observations. The complex flow around the leading edge, resulting in leading-edge separation and consecutive turbulent early reattachment, as witnessed by the pressure distribution, results in a change in the surface temperature distribution at 5%-chord. The turbulent flow separates again at 98%-chord leading to another change in surface temperature gradient. Good agreement is found between IR-thermography and visualizations for the location of the separation point (as in [23]) which is predicted 8%-chord upstream by the CFD. The flow patterns observed in the previous results are summarized in Fig. 7 showing the dependency on the angle of attack of the separation/reattachment points together with the critical point and transition location. On top of the results deduced from the IR-thermography and the CFD, the critical and transition point found using XFOIL [48] are shown. The agreement on the transition location between CFD and IR-thermography is of the order of 5%-chord as is also reported in [14], [23], [46], [49] for other airfoil-Mach combinations. This agreement is met even in the presence of a laminar separation bubble. The discrepancies with XFOIL are of the order of 10%-chord as reported in [46]. XFOIL does not predict the occurrence of a LSB between 3° and 5° angle-of-attack.

The overall agreement pertaining to the separation/reattachment points is of the order of 5%-chord. The distinct flow patterns are in relation with the observed lift and drag coefficients shown in Fig. 8. The figure compares the CFD results with the measured coefficients, those collated in [4] in the incompressible limit, and those obtained with XFOIL [48] (with $N_{crit} = 10$ corresponding to a clean wind tunnel).

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Fig. 8. Lift and drag coefficients for the cambered NACA16009 airfoil (C_{ld} =0.3675)

The measured forces are corrected for solid and wake blockage according to [50]. The qualitative behavior of the results is similar although large discrepancies exist. Both the CFD and the force measurements bear the trace of the distinct flow patterns identified in the previous paragraphs. Particularly, the LSB (3° and 5° angle-ofattack) comes with relatively low drag and high lift when compared to the surrounding angles. The installation of the fully turbulent flow above 5° angle-of-attack comes with a sharp increase in drag and a small drop in lift, thereby substantiating the comments on high angle-ofattack performance of the NACA 16-009 in [4], although the cause of this behavior is not that the flow is fully separated but that it is fully turbulent.

VI. Conclusion

A cambered NACA16-series airfoil (NACA 16-409 with $C_{ld} = 0.3675$) has been investigated at low Mach and Reynolds numbers as would typically be encountered during take-off/landing and initial climb when used in a propeller blade. The results from infrared thermography, force measurements, oil flow visualizations, transitional SST k- ω & γ RANS simulations, and XFOIL are in overall good agreement and depict three distinct flow patterns:

- 1. From 0° to less than 3° angle-of-attack: attached laminar flow that separates early (75%- to 80%-chord), the separated shear layer becomes turbulent but does not reattach;
- From 3° to less than 5° angle-of-attack: a laminar separation bubble extends from about 73%- to 93%-chord, the separated shear layer is transitional at 84%-chord, the reattached turbulent flow separates again 2% up-stream of the trailing edge;
- 3. Above 5° to 9° angle-of-attack: attached turbulent

flow but for the existence of a leading edge separation bubble, the turbulent boundary layer separates around 90%-chord.

In light of these patterns, the reported high angle-ofattack efficiency loss of the NACA 16-series family can be better understood. Next, we believe the present investigations also substantiate the following remarks:

- The use of infrared thermography, and in particular the analysis of low pass filtered temperature profiles for the presence of inflexion points, provides interesting perspectives. However, the technique cannot discriminate for late reattachment and the use of a turbulent witness-flow is mandatory;
- The local correlation-based transition model is able to predict laminar separation bubbles as well as simple separation phenomena with the same accuracy as eddy-viscosity-based phenomenological transition models;
- Although XFOIL is a great tool for the prediction of lift and drag characteristics of airfoils, the reality of the flow behind these characteristics is quite different from the predicted behavior in particular for the NACA16-series and its sharp leading edge. This might explain part of the discrepancies observed in propeller applications when using XFOIL for blade element analysis.

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